

Gamma ray production resulting from the annihilation of neutrino/antineutrino emitted from the accretion disk surrounding compact objects

Zoltán Kovács  
The University of Hong Kong

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# Overview

- Toy model for (anti)neutrino annihilation in accreting systems (disk+central object)
- EOS of the central object (neutron/quark stars)
- Energy production close to the disk
- Energy production along the rotational axis
- Summary & outlook

# Toy model for the energy source

- Rotating compact object  
EOS tables for neutron and quark matter
- Standard accretion disk model  
steady state model, geometrically thin and optically thick disk
- Neutrino source: only the disk, star is neglected
- electron/positron pair creation and  $E$  liberation via neutrino/antineutrino annihilation
- Considered along the equatorial plane and the rotational axis

# EOS

Neutron stars:

DH (Douchin & Hanselm 2001),

RMF soft/stiff (Kubis & Kutchera 1997),

STOS  $T=0, 0.5, 1$  MeV,

(Shen, Toki, Oyamatsu & Sumiyoshi 1998),

BBBAV14 & BBBParis

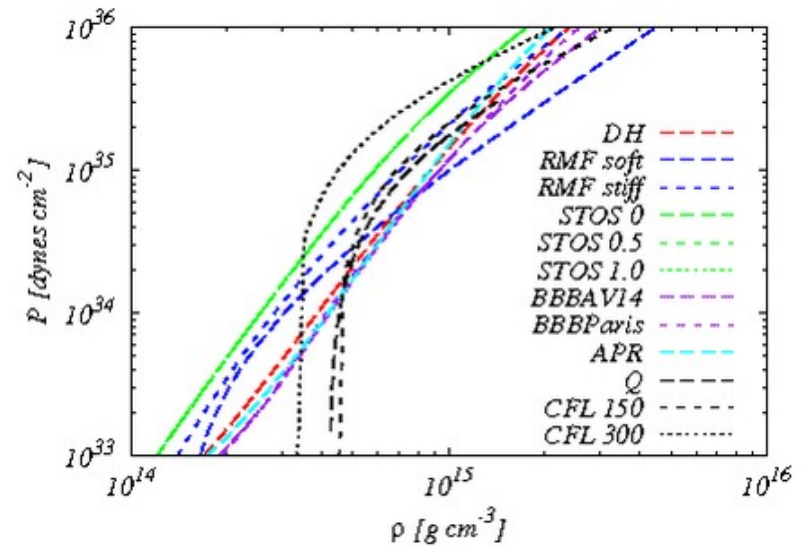
(Baldo, Bombaci & Burgio 1997),

APR (Akmal, Pandharipande & Ravenhall 1998)

Quark stars:

Q (Witten 1984, Chen et al. 1998),

CFL  $\Delta=150$  MeV, ... (Alford et al. 1999)



# Energy Deposition Rate (EDR)

The EDR per unit volume:

$$\dot{q}(\mathbf{r}) = \frac{dE_0}{dt dV} = \int \int f_\nu(\mathbf{p}_\nu, \mathbf{r}) f_{\bar{\nu}}(\mathbf{p}_{\bar{\nu}}, \mathbf{r}) \{ \sigma_{\nu\bar{\nu}} |\mathbf{v}_\nu - \mathbf{v}_{\bar{\nu}}| \varepsilon_\nu \varepsilon_{\bar{\nu}} \} \frac{\varepsilon_\nu + \varepsilon_{\bar{\nu}}}{\varepsilon_\nu \varepsilon_{\bar{\nu}}} d^3 \mathbf{p}_\nu d^3 \mathbf{p}_{\bar{\nu}},$$

$$\sigma_{\nu\bar{\nu}} = KG_F^2 (\varepsilon_\nu \varepsilon_{\bar{\nu}} - c^2 \mathbf{p}_\nu \cdot \mathbf{p}_{\bar{\nu}}) = -KG_F^2 \mathbf{p}_\nu \cdot \mathbf{p}_{\bar{\nu}}$$

Integrating in spherically symmetric geometry of the ST:

$$\begin{aligned} \dot{q}(r) &= \frac{dE_0}{dt dV} = 2cKG_F^2 \Theta(r) \int \int f_\nu f_{\bar{\nu}} (\varepsilon_\nu + \varepsilon_{\bar{\nu}}) \varepsilon_\nu^3 \varepsilon_{\bar{\nu}}^3 d\varepsilon_\nu d\varepsilon_{\bar{\nu}} \\ &= \frac{21\pi^4}{4} \zeta(5) \frac{KG_F^2}{h^6 c^5} k^9 T_{eff}^9 (3r_g) \Theta(r) \end{aligned}$$

# EDR along the disk

Integrated over the 3-volume (Salmonson & Wilson 1999):

$$\dot{Q}(R) = 2 \int_0^{2\pi} \int_0^{\pi/2} \int_R^\infty \dot{q}(r, R, \theta) \sqrt{g_{rr} g_{\theta\theta} g_{\phi\phi}} dr d\theta d\phi.$$

Restricted into the equatorial plane:

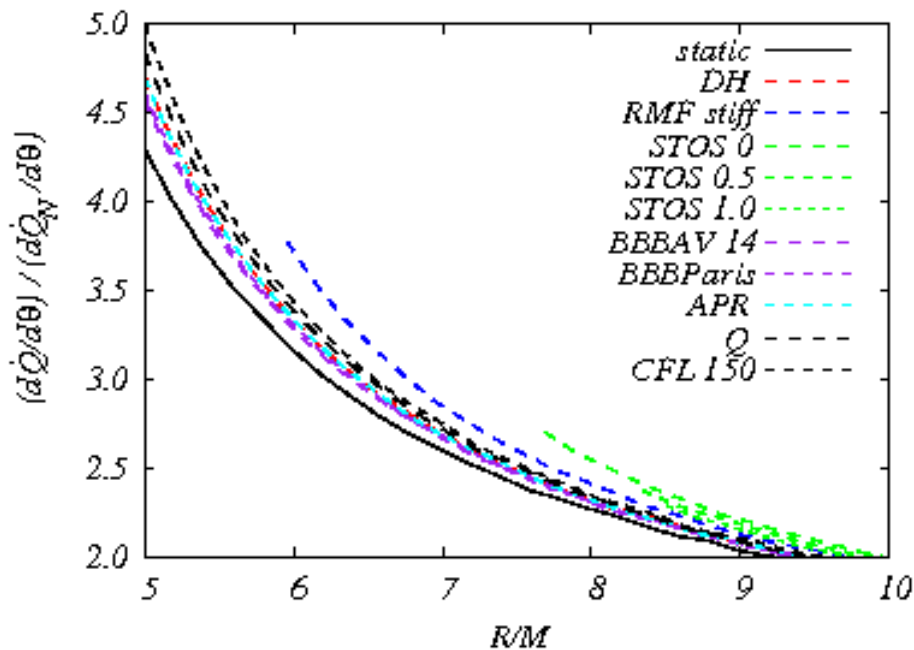
$$\begin{aligned} \left. \frac{d\dot{Q}}{d\theta} \right|_{\theta=\pi/2} &= 2 \int_0^{2\pi} \int_R^\infty \dot{q}(r, R, \theta) \sqrt{g_{rr} g_{\theta\theta} g_{\phi\phi}} dr d\phi \Big|_{\theta=\pi/2} \\ &= 4\pi \int_R^\infty \dot{q}(r, R, \theta) \sqrt{g_{rr} g_{\theta\theta} g_{\phi\phi}} dr \Big|_{\theta=\pi/2}. \end{aligned}$$

Compared with the Newtonian model:

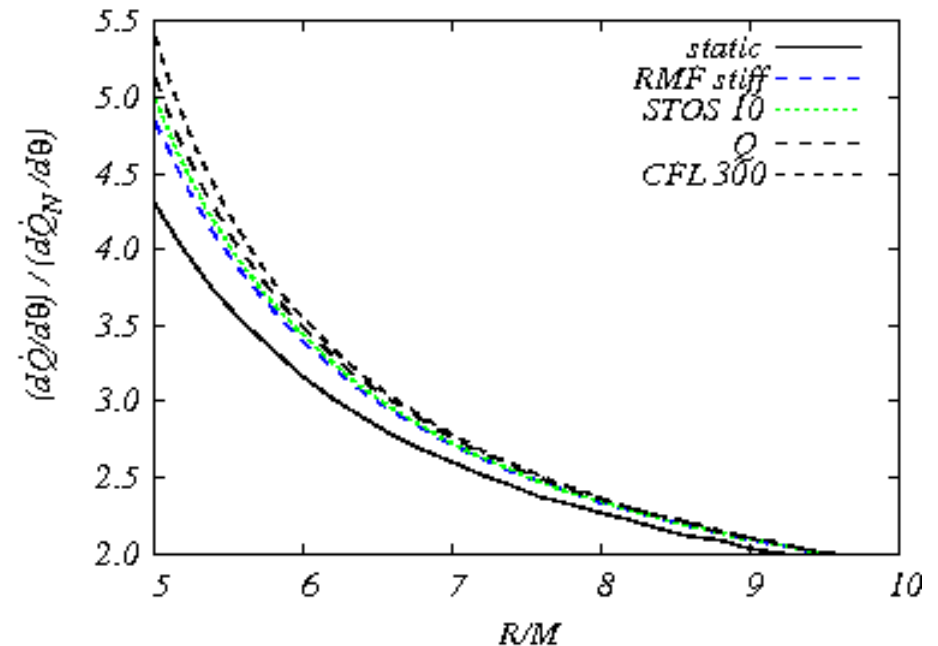
$$\left. \frac{d\dot{Q}/d\theta}{d\dot{Q}_N/d\theta} \right|_{\theta=\pi/2} = \frac{\int_R^\infty \dot{q}(r, R, \theta) \sqrt{g_{rr}(r, \theta) g_{\theta\theta}(r, \theta) g_{\phi\phi}(r, \theta)} dr \Big|_{\theta=\pi/2}}{\int_R^\infty \dot{q}_N(r, R) r^2 dr}.$$

# EDR along the disk

The radial distribution of EDR restricted into the equatorial plane (Kovács, Cheng & Harko 2009)



$$M = 1.8 M_{\odot}, \quad \Omega = 5 \times 10^3 \text{ s}^{-1}$$

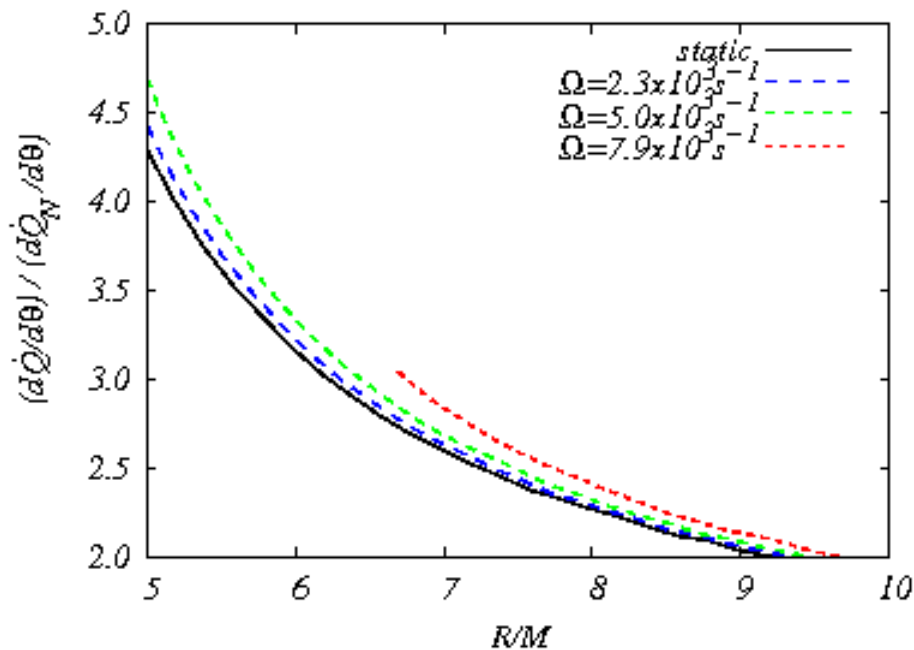


$$M = 2.8 M_{\odot}, \quad \Omega = 5 \times 10^3 \text{ s}^{-1}$$

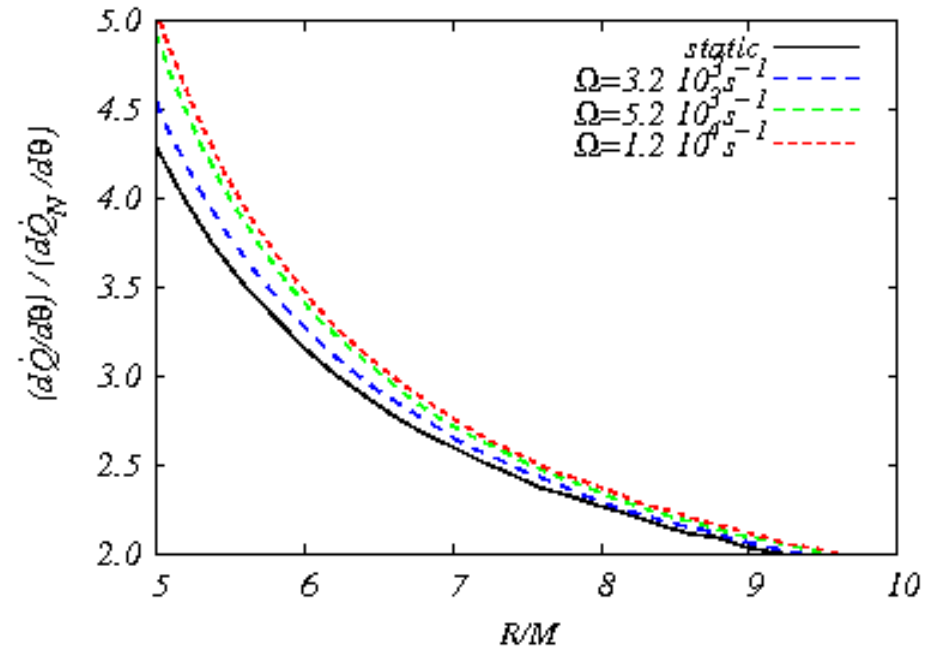
RMF stiff & STOS: high EDR but small disk surface.  
Quark stars produce higher EDR.

# EDR along the disk

Dependence of the EDR on  $\Omega$  for APR and Q type EOS  
(Kovács, Cheng & Harko 2009)



$M = 1.8 M_{\odot}$ , APR



$M = 1.8 M_{\odot}$ , Q

The EDR is proportional to the rotational frequency.



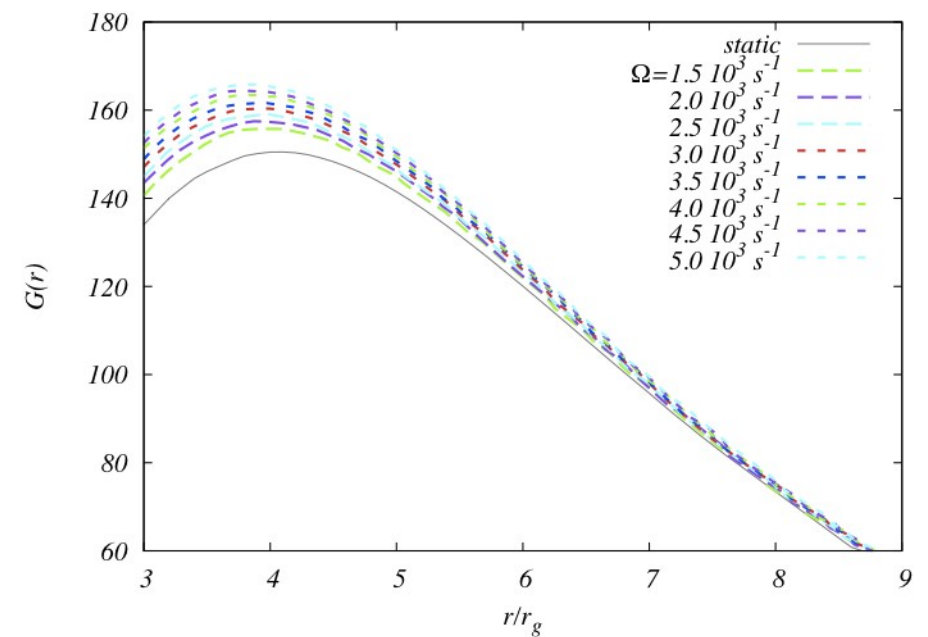
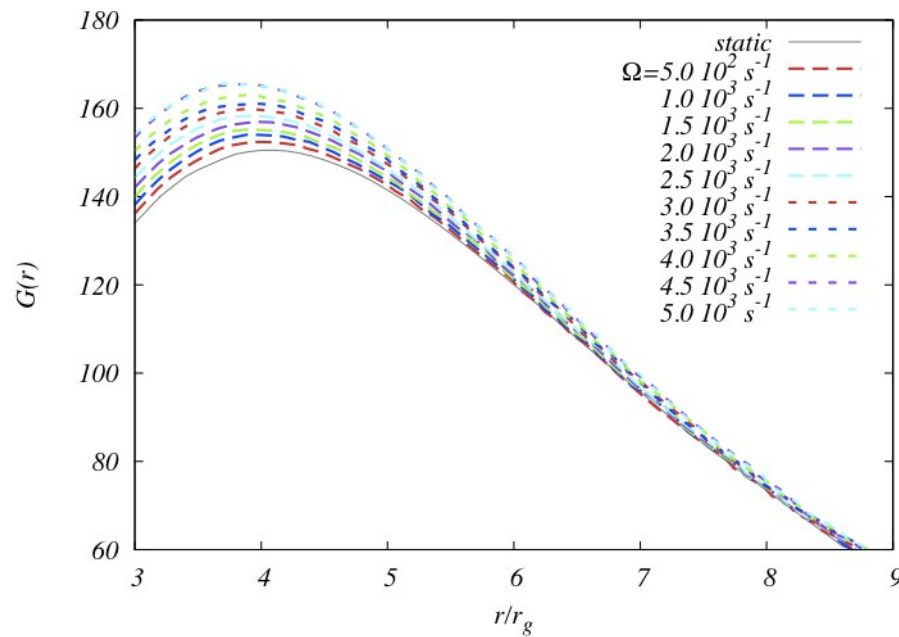
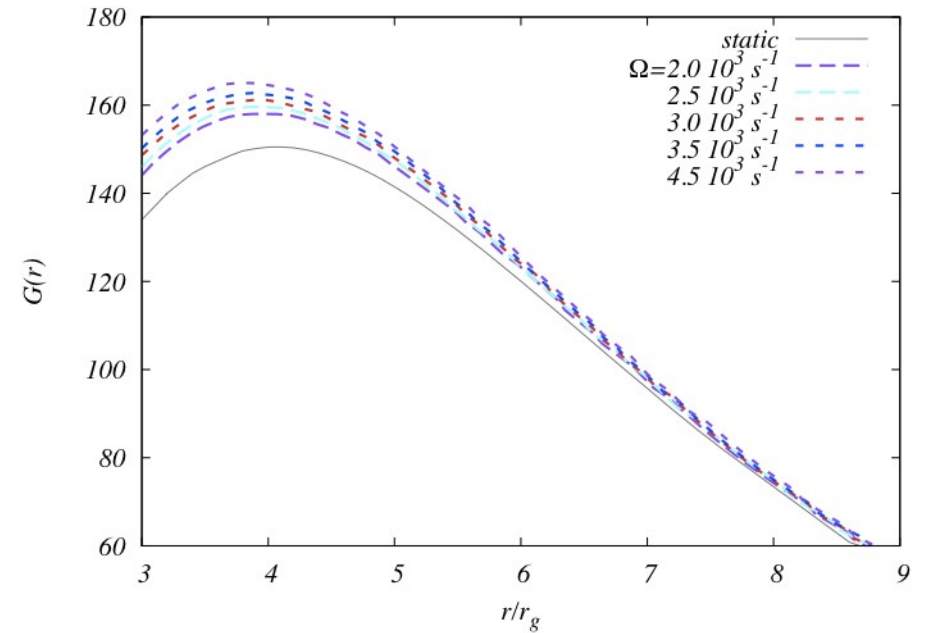
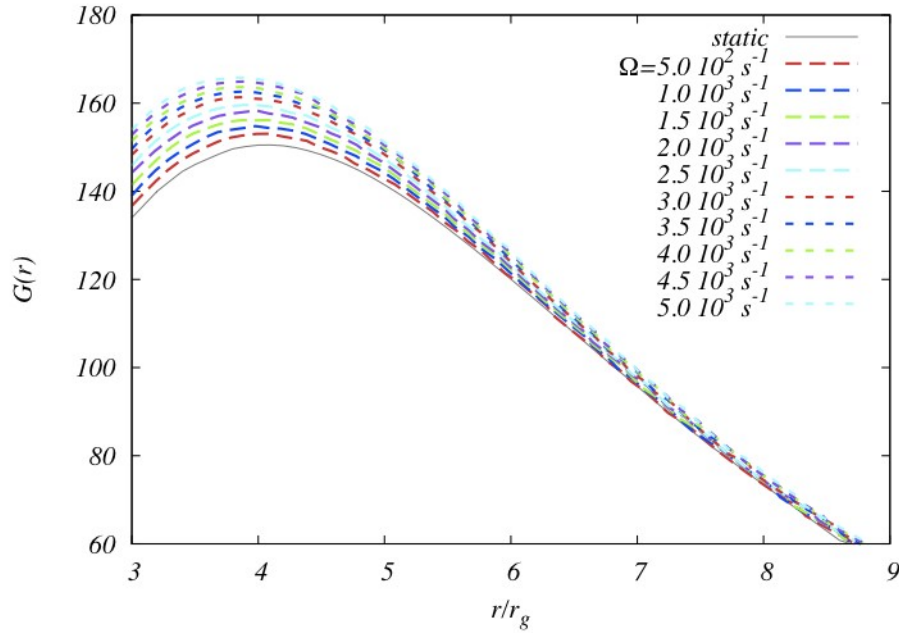
# EDR along the rotational axis

Integrated over the 4-volume (Asano & Fukuyama 2001) but restricted along the axis of rotation:

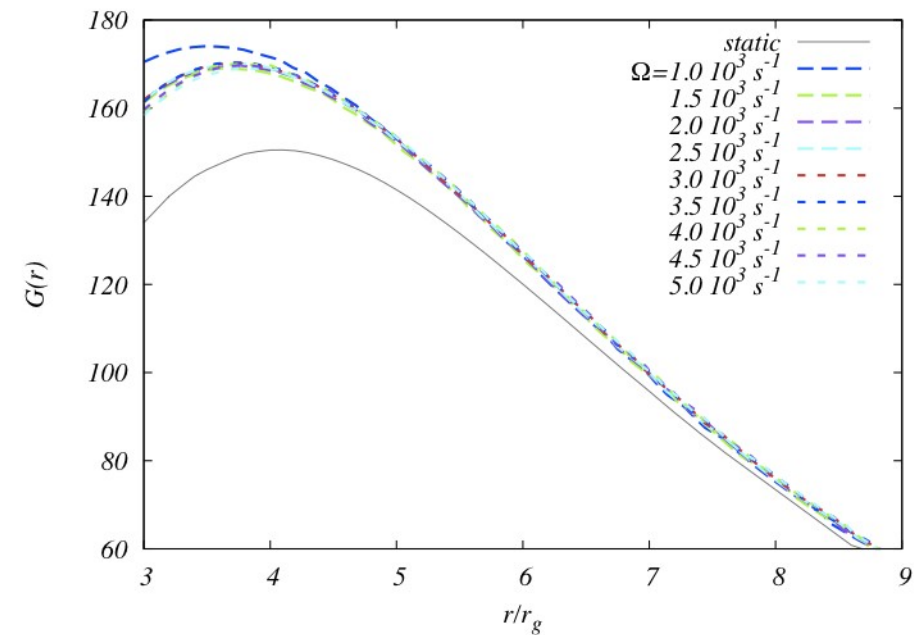
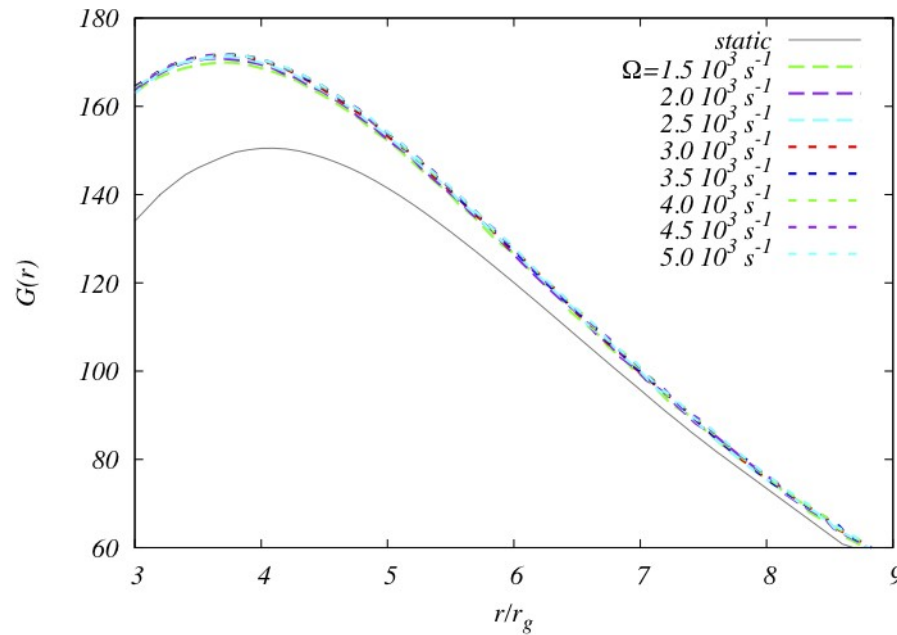
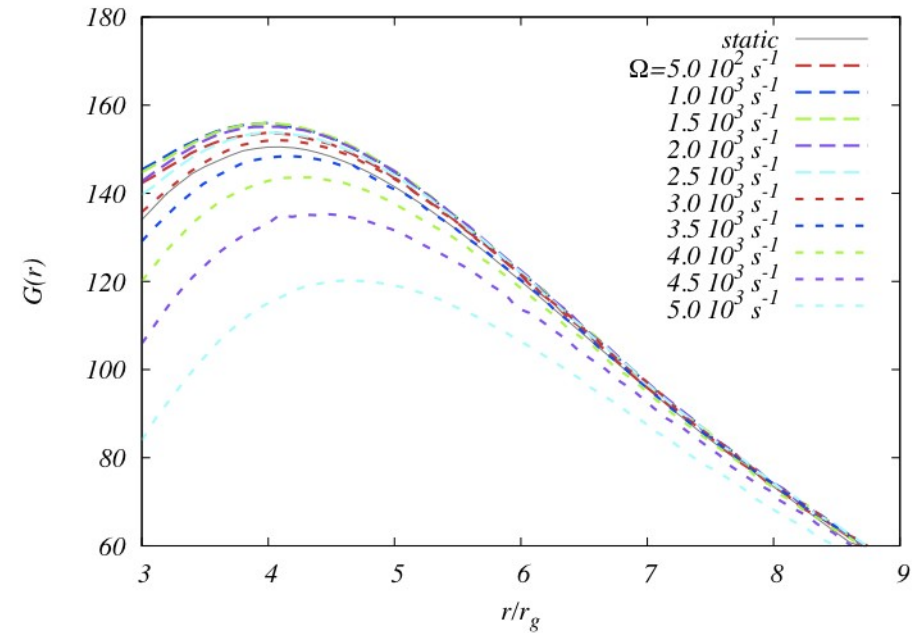
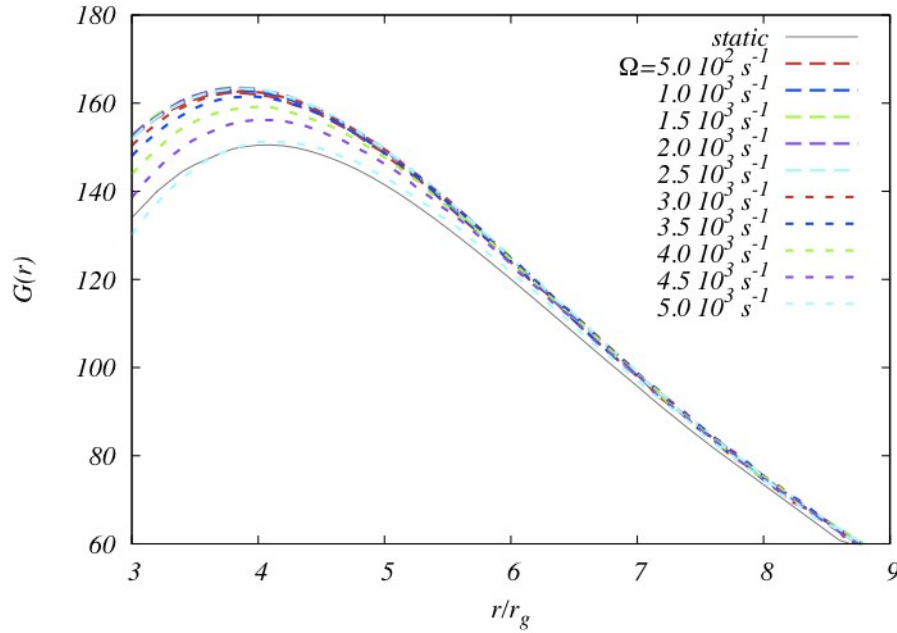
$$\begin{aligned}\left. \frac{dE_0}{dt} \right|_{\theta=0} &\simeq 2 \frac{d}{d\theta} \int_0^{2\pi} \int_0^{\pi/2} \int_{r_{min}}^{r_{max}} \left. \frac{dE_0}{dt dV} \sqrt{-g} dr d\theta d\phi \right|_{\theta=0} \\ &= 21\pi^5 \zeta(5) \frac{K G_F^2}{h^6 c^5} k^9 T_{eff}^9 (3r_g) r_g^2 \int_{r_{min}}^{r_{max}} G(r) dr , \\ G(r) &\equiv r_g^{-2} \left. \frac{\sqrt{-g(r, \theta)}}{\sin \theta} \right|_{\theta=0} \Theta(r) .\end{aligned}$$

G describes the effects of the geometry on the EDR.

# Dependence of the EDR on $\Omega$ : DH, APR, BBBAV14 & BBBParis



# Dependence of the EDR on $\Omega$ : RMF stiff, STOS, Q & CFL



# Summary & Outlook

- Studied simple models for energy production due to antineutrino/neutrino annihilation in accreting systems with rotating central objects
- Considered a broad variation of EOS for the central neutron and quark stars
- Presented the radial distribution of EDR along the equatorial plane and the rotational axis
- Considerable dependence on the EOS: quark stars produce higher rates
- Possible reconstruction of 3D maps for EDR