Computer simulations of strongly correlated quantum systems: past successes, current topics, and future challenges

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with a guest appearance by the elusive supersolid

Phases and phase transitions



Numerical simulations

"You let the computer solve the problem for you"

- □ It's not that easy:
 - Exponentially diverging number of states
 - Critical slowing down of the dynamics at phase transitions





Х

Ulam: the Monte Carlo Method

□ What is the probability to win in Solitaire?

Ulam's answer: play it 100 times, count the number of wins and you have a pretty good estimate





The Monte Carlo Method



Need a representative sample with the correct distribution

$$P[c_i] = \frac{p_{c_i}}{Z}$$

fundamental problem of statistical mechanics

- Metropolis Algorithm for Monte Carlo
- Simplex Method for Linear Programming

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- Krylov Subspace Iteration Methods
- The Decompositional Approach to Matrix Computations
- The Fortran Optimizing Compiler
- QR Algorithm for Computing Eigenvalues
- Quicksort Algorithm for Sorting
- Fast Fourier Transform
- Integer Relation Detection
- Fast Multipole Method



The Metropolis Algorithm (1953)

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Equation of State Calculations by Fast Computing Machines

NICHOLAS METROPOLIS, ARIANNA W. ROSENBLUTH, MARSHALL N. ROSENBLUTH, AND AUGUSTA H. TELLER, Los Alamos Scientific Laboratory, Los Alamos, New Mexico

AND

EDWARD TELLER,* Department of Physics, University of Chicago, Chicago, Illinois (Received March 6, 1953)

A general method, suitable for fast computing machines, for investigating such properties as equations of state for substances consisting of interacting individual molecules is described. The method consists of a modified Monte Carlo integration over configuration space. Results for the two-dimensional rigid-sphere system have been obtained on the Los Alamos MANIAC and are presented here. These results are compared to the free volume equation of state and to a four-term virial coefficient expansion.

I. INTRODUCTION

THE purpose of this paper is to describe a general method, suitable for fast electronic computing machines, of calculating the properties of any substance which may be considered as composed of interacting individual molecules. Classical statistics is assumed,

II. THE GENERAL METHOD FOR AN ARBITRARY POTENTIAL BETWEEN THE PARTICLES

In order to reduce the problem to a feasible size for numerical work, we can, of course, consider only a finite number of particles. This number N may be as high as several hundred. Our system consists of a square[†] con-

I. INTRODUCTION

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The Metropolis Algorithm (1953)

creates a representative sample for any system start with a configuration i $\uparrow \uparrow \uparrow \uparrow \uparrow \uparrow \uparrow \uparrow$ propose a small change to a configuration j $\uparrow \uparrow \uparrow \uparrow \downarrow \uparrow \uparrow \uparrow$

calculate the ratio of weights $\frac{p_j}{p_i}$

accept the new configuration with probability

$$P = \min\left(1, \frac{p_j}{p_i}\right)$$

Quantum Monte Carlo

Feynman (1953) lays foundation for quantum Monte Carlo
 Map quantum system to classical world lines

PHYSICAL REVIEW

ТНЕ

A journal of experimental and theoretical physics established by E. L. Nichols in 1893

Second Series, Vol. 91, No. 6

SEPTEMBER 15, 1953

Atomic Theory of the α Transition in Helium

R. P. FEYNMAN California Institute of Technology, Pasadena, California (Received May 15, 1953)

Use Metropolis algorithm to update world lines

Autocorrelation effects

The Metropolis algorithm creates a Markov chain of configurations

$$c_1 \to c_2 \to \dots \to c_i \to c_{i+1} \to \dots$$

successive configurations are correlated, leading to an increased statistical error

$$\Delta A = \sqrt{\left\langle \left(\overline{A} - \left\langle A \right\rangle\right)^2 \right\rangle} = \sqrt{\frac{\operatorname{Var} A}{M}} (1 + 2\tau_A)$$

Critical slowing down at second order phase transition

 $\tau \propto L^2$

Exponential tunneling problem at first order phase transition $\tau \propto \exp(L^{d-1})$

Modern algorithms

Critical slowing down at a 2nd order phase transition is solved by *changing the dynamics*: make large global changes
 Swendsen and Wang, 1987; Evertz *et al*, 1993; Prokof'ev *et al*, 1998 ...

 Tunneling problem at a first order phase transition is solved by *changing the ensemble* to create a flat energy landscape
 Berg & Neuhaus, 1992; Wang & Landau, 2001; Troyer *et al*, 2003, ...



Improvements in simulation methods

 30 years of simulating the Ising model by D.P. Landau
 Today's algorithms on 30 year old computers faster than 30 year old algorithms on today's computers



Modern algorithms for quantum systems

 Larger systems at lower temperature allow accurate simulation of phase transitions in quantum systems

Temperature	Metropolis	Modern algorithms
J	16'000 spins	16'000'000 spins
0.1 J	200 spins	1'000'000 spins
0.005 J		50'000 spins
0.I t	32 lattice bosons	10'000 lattice bosons
50 mK	800 He atoms	8'000 He atoms

Experiments on the computer

Modern QMC algorithms allow accurate simulation of quantum many body problems: experiments on models

Example: the search for the supersolid



lattice supersolids

is solid Helium a supersolid?

Supersolid

- A supersolid is simultaneously
 - superfluid: broken gauge symmetry
 solid: broken translational symmetry
- First proposals nearly as old as I am
 A. F. Andreev and I. M. Lifshitz, Sov. Phys. JETP 29, 1107 (1969).
 G. V. Chester, Phys. Rev. A 2, 256 (1970).

but not seen in Nature yet

Super solids in lattice boson models

The simplest proposals are lattice models:
Square lattice hardcore bosons with repulsion

$$H = -t \sum_{\langle i,j \rangle} (a_i^{\dagger} a_j + a_j^{\dagger} a_i) + V \sum_{\langle i,j \rangle} n_i n_j$$



solid (crystal) caused by large V

doped solid: interstitials supersolid: superfluid interstitials

First simulation results: it exists!

Batrouni, Scalettar, Zimanyi, Kampf, PRL (1995)

- evidence for supersolid at 3% doping
- finite superfluid density and solid structure factor



Do lattice supersolids exist?

D Bosons with nearest neighbor repulsion

- mean-field calculations found supersolid
- **previous simulations** (32 particles) found **supersolid**
- new simulations (5000 particles) instead show phase separation at first order phase transition





Supersolids versus phase separation



solid

doped solid







supersolid



doped particles gain energy by forming a domain wall

Stabilizing the supersolid

add next nearest neighbor hopping: $\tilde{H} = H - t' \sum_{\langle \langle i,j \rangle \rangle} (a_i^{\dagger} a_j + a_j^{\dagger} a_i)$ Sengupta et al, PRL 2005







form a striped solid by longer range repulsion





Batrouni *et al*, PRB 1995 Batrouni and Scalettar, PRL 2000 Hebert *et al*, PRB 2002 Schmid *et al*, PRL 2002 Schmid and Troyer, PRL 2004



Wessel and Troyer, PRL 2005 Melko *et al*, PRL 2005 Heidarian and Damle, PRL 2005 Boninsegni *et al*, PRL 2005

Continuum supersolids

Supersolidity due to superflow of vacancies in solid
 A. F. Andreev and I. M. Lifshitz, Sov. Phys. JETP 29, 1107 (1969).
 G. V. Chester, Phys. Rev. A 2, 256 (1970).



A related scenario proposed recently by
 X. Dai, M. Ma, F.-C. Zhang, Phys. Rev. B 72, 132504 (2005)
 P.W. Anderson, W.F. Brinkman and D.A. Huse, Science 310, 1164 (2005).

Experiment by Kim and Chan



Reproduced by

- A. S. Rittner and J. D. Reppy, PRL (2006)
- M. Kubota et al., unpublished (2006)
- K. Shirahama et al., unpublished (2006).

E. Kim and M.H.W. Chan nonclassical moment of inertia Nature, **427**, 225 (2004) Science **305**, 1941 (2004)



Vacancies and interstitials are gapped

Energy obtained from decay of Green's function



M. Boninsegni, N. V. Prokof'ev and B. V. Svistunov, PRL 2006 B. K. Clark and D. M. Ceperley, PRL 2006 M. Boninsegni, A. Kuklov, L. Pollet, N.V. Prokof'ev, B.V. Svistunov, M. Troyer, PRL 2006

Theories : overview



Phase separation of vacancies

Strong attraction between vacancies

Vacancies phase separate: solid Helium is not supersolid



Quantum: M. Boninsegni, A. Kuklov, L. Pollet, N.V. Prokof'ev, B.V. Svistunov, M. Troyer, PRL 2006 Classical: Tama Ma (HKU), L. Pollet and M. Troyer, in preparation

Instability of the supersolid

Theory of supersolid is based on dilute gas of bosonic vacancies
 stability of a Bose-Einstein condensate requires repulsive interactions
 unstable in Helium due to attractive interaction between vacancies

□ Future challenges:

- what is "wrong" with the Helium potential?
- which potential could give a supersolid?
- what is seen in the Kim & Chan experiment?
- maybe due to superflid grain boundaries? (simulation and experiment)

Simulation of grain boundary

Points on exchange cycles indicate superfluid regions

insulating solid 3 5 T 49 insulating solid superfluid grain boundary

superfluid

superfluid

Future challenges for simulations

Dynamics of quantum systems

- Monte Carlo simulations perfect for static equilibrium properties
- Dynamic properties and non-equilibrium effects are hard to obtain

Entropy-driven phase transitions

- It is hard to form a single crystal from a liquid in experiment
- Even harder to form a single crystal from a liquid in simulations

Fermionic simulations

The infamous negative sign problem prevents accurate simulations of fermions

The negative sign problem

In mapping of quantum to classical system

$$Z = \mathrm{Tr}e^{-\beta H} = \sum_{i} p_i$$

 \square there is a "sign problem" if some of the $p_i < 0$

Appears e.g. in simulation of fermions when two fermions exchange places, since the sign of the wave function changes (Pauli principle)



The negative sign problem

Sample with respect to absolute values of the weights

$$\langle A \rangle = \sum_{i} A_{i} p_{i} / \sum_{i} p_{i} = \frac{\sum_{i} A_{i} \operatorname{sgn} p_{i} |p_{i}| / \sum_{i} |p_{i}|}{\sum_{i} \operatorname{sgn} p_{i} |p_{i}| / \sum_{i} |p_{i}|} \equiv \frac{\langle A \cdot \operatorname{sign} \rangle_{|p|}}{\langle \operatorname{sign} \rangle_{|p|}}$$

Exponentially growing cancellation in the sign

$$\langle sign \rangle = \frac{\sum_i p_i}{\sum_i |p_i|} = Z/Z_{|p|} = e^{-\beta V(f - f_{|p|})}$$

Exponential growth of errors

$$\frac{\Delta sign}{\langle sign \rangle} = \frac{\sqrt{\langle sign^2 \rangle - \langle sign \rangle^2}}{\sqrt{M} \langle sign \rangle} \approx \frac{e^{\beta V (f - f_{|p|})}}{\sqrt{M}}$$

□ NP-hard problem (no general solution) [Troyer and Wiese, PRL 2005]

How to overcome the sign problem

□ Full solution is impossible

- Exact solution in one dimension and small clusters
 - exact diagonalization of up to about 30-40 sites
 - density matrix renormalization group method on 1000 sites in 1D

Approximate solutions in higher dimensions

- Dynamical mean field theory embeds a small cluster in a selfconsistent bath: can model the Mott-transition well
- Variational Monte Carlo methods, such as fixed-node approximation, can work well but it is hard to unambiguously establish the existence of new exotic phases

Improved algorithms and new ideas are needed

Example: new DMFT solvers

Substantial progress is still possible

- Standard method: Hirsch & Fye, PRL '86
- New weak coupling expansion: Rubtsov et al, PRB '05
- New hybridizaton expansion: Werner et al, PRL '06



■ New methods allow

- 100 x lower temperatures
- accurate investigation of Mott transition
- and more to come ...

Computational Physics



Computational Physics has become an important branch of physics, complementary to experiment and theory Perform "experiments on a model"

Summary

- Modern algorithms allow accurate simulations of phase diagrams and phase transitions of many quantum models
- But substantial challenges remain
 - dynamics
 - entropy-driven phase transitions
 - fermionic simulations
 - Example: the search for a supersolid
 - lattice supersolids can be realized using ultracold polar molecules in triangular optical lattices
 - Helium-4 is according to our simulations not a supersolid: defects play an important role in experiments

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